

ON THE GEOMETRY OF THE CALOGERO-MOSER SYSTEM

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To Gerrit van Dijk for his 65th birthday

ABSTRACT. We discuss a special eigenstate of the quantized periodic Calogero-Moser system associated to a root system. This state has the property that its eigenfunctions, when regarded as multivalued functions on the space of regular conjugacy classes in the corresponding semisimple complex Lie group, transform under monodromy according to the complex reflection representation of the affine Hecke algebra. We show that this endows the space of conjugacy classes in question with a projective structure. For a certain parameter range this projective structure underlies a complex hyperbolic structure. If in addition a Schwarz type of integrality condition is satisfied, then it even has the structure of a ball quotient minus a Heegner divisor. For example, the case of the root system E_8 with the triflection monodromy representation describes a special eigenstate for the system of 12 unordered points on the projective line under a particular constraint.

1. INTRODUCTION

The classical Calogero-Moser system describes a finite number of identical point particles on the real line under the influence of an inverse square potential. The relative positions of these points are parametrized by the complement of a hyperplane arrangement, namely by \mathbb{R}^{n+1} modulo its main diagonal minus the union of the hyperplanes for which two coordinates are equal (in case they are $n + 1$ in number). Thus we find ourselves immediately dealing with a root system of type A_n . It is therefore not surprising that there is a generalization to arbitrary finite root systems. This can be (and has been) taken a step further by letting the basic input be an abstract root system R and a one dimensional connected complex Lie group \mathbb{G} : if P denotes the weight lattice of R , then the substitute for the configuration space is the abelian complex Lie group $\text{Hom}(P, \mathbb{G})$ deprived from the fixed point divisors of the reflections in the Weyl group W of R . So if \mathbb{G} is the additive group \mathbb{C} , then we recover the previous generalization, but if \mathbb{G} is the multiplicative group \mathbb{C}^\times resp. an elliptic curve, then we get the complement of a reflection arrangement in an algebraic torus resp. an abelian variety. It is common terminology to talk about the non-periodic, the periodic and the doubly periodic case of the root system generalization of the Calogero-Moser system.

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The general integrability properties of these systems are nowadays rather well understood. Moser showed that the classical (A_n) system is completely integrable in the additive and multiplicative case by realizing such a system as a Lax pair [19]. Olshanetsky and Perelomov extended his method to the elliptic curve case [20] and also began the study of the Calogero-Moser system in the general context of root systems [21],[22]. The question of complete integrability for the quantization of the Calogero-Moser system in the multiplicative case and for a general root system R was solved by complex analytic methods in the work of Heckman and Opdam [13]. The associated eigenvalue system is also called the hypergeometric system associated with R because in the rank one situation it boils down to the Euler-Gauss hypergeometric equation. Heckman later showed that these results could be obtained in an elementary way by means of the theory of Dunkl operators [10]. Subsequently Cherednik and Matsuo noticed the analogy between this approach and the Knizhnik-Zamolodchikov equation of conformal field theory [6], [18]. Cherednik also made the important observation that the relations for the (trigonometric) Dunkl operators are those of the degenerate affine Hecke algebra of Drinfeld and Lusztig [9], [23], [11]. The affine Hecke algebra appeared in the monodromy of the hypergeometric system associated with R as well [13], and the role of the (double) affine Hecke algebra in the basic hypergeometric extension (where q -difference operators replace differential operators) became increasingly important [17].

It turns out that our previous work on the geometric structures on complements of projectivized hyperplane arrangements can be understood in and generalized to this context. Our paper [7] accomplished that in the additive setting by generalizing the work of Deligne-Mostow on the Lauricella functions (which corresponds from our perspective to the classical A_n -case) to arbitrary root systems.

In the present article we describe the multiplicative version, leaving the case of an elliptic curve for another occasion. Two sets of free parameters enter here: the system of linear differential equations depends on a set of *coupling parameters* and their common eigen space decomposition is indexed by *spectral parameters*. We focus on the situation when there is a subsystem of the hypergeometric system that is associated to the *reflection representation* of the affine Hecke algebra; this amounts to a particular choice of the spectral parameters. This turns out to define a projective structure on the space of W -orbits of the toric arrangement complement in question. For a certain range of the coupling parameters this even yields a complex hyperbolic structure on that orbit space. It is usually incomplete, but we find that when, in the terminology of [7], the coupling parameters satisfy the *Schwarz conditions*, the situation is as nice as one could hope for: we thus obtain a handful of cases for which the hyperbolic structure can be completed to a quotient of a complex ball by a discrete group of automorphisms acting with cofinite invariant volume. We sketch these results in Section 3 with the intention of treating this material more thoroughly elsewhere.

We found it worthwhile to discuss in some detail (in the final Section 4) the example where R is of type E_8 and the coupling parameter equals $\frac{1}{6}$, because its

geometry is linked to the moduli spaces of rational elliptic surfaces and degree 12 divisors on the projective line [12]. Similar stories can be told for E_6 with coupling parameter $\frac{1}{3}$ and E_7 with coupling parameter $\frac{1}{4}$ in relation to moduli spaces of cubic surfaces and quartic curves respectively [1], [15]. But despite these beautiful examples, the picture of the geometry behind the general Calogero-Moser system (with general R and arbitrary spectral and coupling parameters) remains unclear. Recent papers on the classical Calogero-Moser system point in the direction of the classical Hitchin system [14], [4]. A natural continuation along these lines seems to lead to a further analysis of the quantization of Hitchin's fibration [2], [3].

For the reader's convenience we included a brief review of the theory of trigonometric Dunkl operators in Section 2.

We dedicate this paper to Gerrit van Dijk on the occasion of his 65th birthday. One of us (GH) was the first PhD student of Gerrit van Dijk at Leiden University, and as such he is grateful for the excellent mathematical training he received from Gerrit.

2. QUANTIZATION OF THE TRIGONOMETRIC CALOGERO-MOSER SYSTEM

Let \mathfrak{a} be a real vector space of dimension n , and let $R \subset \mathfrak{a}^*$ be a possibly nonreduced root system spanning \mathfrak{a}^* . Let $R^\vee \subset \mathfrak{a}$ be the dual root system and write $Q^\vee = \mathbb{Z}R^\vee \subset \mathfrak{a}$ for the coroot lattice of R so that $P := \text{Hom}(Q^\vee, \mathbb{Z}) \subset \mathfrak{a}^*$ is the weight lattice of R . Then $H = \text{Hom}(P, \mathbb{C}^\times)$ is a complex torus whose character lattice $\text{Hom}(H, \mathbb{C}^\times)$ can be identified with P and whose Lie algebra $\text{Lie}(H) = \mathfrak{h}$ can be identified with $\mathfrak{a}_\mathbb{C}$. Notice that the algebra $\mathbb{C}[H]$ of regular functions on H can be identified with the group algebra of P , because it has $\{e^\mu\}_{\mu \in P}$ as an additive basis and the multiplication law is simply $e^\mu e^\nu = e^{\mu+\nu}$.

For $\alpha \in R$, we denote by $s_\alpha : \lambda \mapsto \lambda - \lambda(\alpha^\vee)\alpha$ the corresponding reflection of \mathfrak{h}^* . These reflections generate the Weyl group $W \subset \text{GL}(\mathfrak{h}^*)$ of R . By duality W also acts on \mathfrak{h} and H . The reflection s_α fixes the hypertorus $H_\alpha = \{h \in H : e^\alpha(h) = 1\}$ pointwise. We put $H^\circ = H - \bigcup_{\alpha \in R} H_\alpha$. It is clear that the algebra $\mathbb{C}[H^\circ]$ of regular functions on H° is obtained from $\mathbb{C}[H]$ by inverting the elements $1 - e^{-\alpha}$, $\alpha \in R$. For an element p of the symmetric algebra of \mathfrak{h} , $\text{Sym}(\mathfrak{h})$, we denote by $\partial(p)$ the corresponding translation invariant differential operator on H . So if we identify $\text{Sym}(\mathfrak{h})$ with the algebra of polynomial functions on \mathfrak{h}^* , $\mathbb{C}[\mathfrak{h}^*]$, then $\partial(p)e^\mu = p(\mu)e^\mu$ for $\mu \in P$. The algebra of differential operators on H° can (as a $\mathbb{C}[H^\circ]$ -module) be identified with $\mathbb{C}[H^\circ] \otimes_{\mathbb{C}} \text{Sym}(\mathfrak{h})$; we shall denote it by $\mathbb{D}[H^\circ]$. Now $\mathbb{C}[H^\circ]$ is a natural left module for $\mathbb{D}[H^\circ]$. The group algebra $\mathbb{C}[W]$ of W also acts on $\mathbb{C}[H^\circ]$ on the left and together they turn $\mathbb{C}[H^\circ]$ into a left module for the semidirect tensor product $\mathbb{D}[H^\circ] \otimes_{\mathbb{C}} \mathbb{C}[W]$ with multiplication law

$$D \otimes w \cdot D' \otimes w' := (DwD'w^{-1}) \otimes ww'.$$

It is easy to see that $\mathbb{D}[H^\circ] \otimes \mathbb{C}[W]$ acts faithfully on $\mathbb{C}[H^\circ]$.

Choose a nonzero W -invariant inner product $(\ , \)$ on \mathfrak{a}^* and denote by $C \in \text{Sym}^2 \mathfrak{h}$ the associated (complexified) quadratic form. Then $\partial(C)$ is the Laplace operator on H characterized by $\partial(C)(e^\mu) = C(\mu)e^\mu$, $\mu \in P$.

Definition 2.1. Let $\mathbf{k} : \alpha \in R \mapsto k_\alpha \in \mathbb{C}$ be a W -invariant map. Then the quantum Hamiltonian of the trigonometric Calogero-Moser system with coupling vector $\mathbf{k} \in \mathbb{C}^R$ is the element of $\mathbb{D}[H^\circ]$ given by

$$H_{\mathbf{k}} = \partial(C) + \frac{1}{2} \sum_{\alpha \in R} \frac{k_\alpha(1 - k_\alpha - 2k_{2\alpha})(\alpha, \alpha)}{(e^{\alpha/2} - e^{-\alpha/2})^2}.$$

If the coupling vector is real, then the quantum Hamiltonian becomes on the compact subtorus $T := \text{Hom}(P, U(1))$ of H the Schrödinger operator for a single particle on T repelled or attracted along the subtori $T_\alpha = T \cap H_\alpha$ by an inverse square distance potential.

Let us fix a set of positive roots R_+ in R . The complete integrability of the Calogero-Moser system is best understood after formal conjugation by

$$\delta_{\mathbf{k}}^{\frac{1}{2}} = \prod_{\alpha > 0} (e^{\alpha/2} - e^{-\alpha/2})^{k_\alpha} = e^{\rho_{\mathbf{k}}} \prod_{\alpha > 0} (1 - e^{-\alpha})^{k_\alpha},$$

where $\rho_{\mathbf{k}} = \frac{1}{2} \sum_{\alpha > 0} k_\alpha \alpha$. We get

$$\delta_{\mathbf{k}}^{-\frac{1}{2}} H_{\mathbf{k}} \delta_{\mathbf{k}}^{\frac{1}{2}} = L_{\mathbf{k}} + (\rho_{\mathbf{k}}, \rho_{\mathbf{k}}),$$

where

$$L_{\mathbf{k}} = \partial(C) + \frac{1}{2} \sum_{\alpha > 0} k_\alpha(\alpha, \alpha) \frac{1 + e^{-\alpha}}{1 - e^{-\alpha}} \partial(\alpha^\vee).$$

Differential operators in $\mathbb{D}[H^\circ]$ commuting with $L_{\mathbf{k}}$ yield after conjugation by $\delta_{\mathbf{k}}^{\frac{1}{2}}$ differential operators in $\mathbb{D}[H^\circ]$ commuting with $H_{\mathbf{k}}$ and vice versa. The advantage of $L_{\mathbf{k}}$ over $H_{\mathbf{k}}$ will become clear later.

Definition 2.2. For $\xi \in \mathfrak{h}$ the trigonometric Dunkl operator $T_{\mathbf{k}}(\xi) \in \mathbb{D}[H^\circ] \otimes_{\mathbb{C}} \mathbb{C}[W]$ with coupling vector \mathbf{k} is

$$T_{\mathbf{k}}(\xi) = (\partial(\xi) - \rho_{\mathbf{k}}(\xi)) \otimes 1 + \sum_{\alpha > 0} \frac{k_\alpha \alpha(\xi)}{1 - e^{-\alpha}} \otimes (1 - s_\alpha).$$

This is a first order operator which preserves the subalgebra $\mathbb{C}[H]$ of $\mathbb{C}[H^\circ]$ because the divided difference operator $(1 - e^{-\alpha})^{-1} \otimes (1 - s_\alpha)$ has that property.

Let us denote by a bar the antilinear involution of $\mathbb{C}[H]$ which sends e^μ to $e^{-\mu}$ ($\mu \in P$). If the coupling vector \mathbf{k} lies in \mathbb{N}^R , then $\delta_{\mathbf{k}}^{1/2} \overline{\delta_{\mathbf{k}}^{1/2}}$ is a well-defined element of $\mathbb{C}[H]$ and we have a hermitian inner product $(\ , \)_{\mathbf{k}}$ on $\mathbb{C}[H]$ defined by

$$(f, g)_{\mathbf{k}} := (\text{constant term of } f \overline{g} \delta_{\mathbf{k}}^{\frac{1}{2}} \overline{\delta_{\mathbf{k}}^{\frac{1}{2}}}) / |W|.$$

It is easy to check that the Dunkl operators are hermitian with respect to $(\ , \)_{\mathbf{k}}$:

$$(T_{\mathbf{k}}(\xi)f, g)_{\mathbf{k}} = \left(f, \overline{T_{\mathbf{k}}(\xi)g} \right)_{\mathbf{k}} = (f, T_{\mathbf{k}}(\overline{\xi})g)_{\mathbf{k}},$$

where the last bar is complex conjugation on \mathfrak{h} relative to the real form \mathfrak{a} of \mathfrak{h} .

Recall that the usual partial ordering \leq on \mathfrak{h}^* is defined by $\mu \leq \nu$ if and only if $\nu - \mu \in \mathbb{N}R_+$. Define a new partial ordering \leq_+ on P as follows. Let P_+ be the

cone of dominant weights. If for $\mu \in P$, μ_+ denotes the unique dominant weight in the orbit $W\mu$, then we let

$$\mu \leq_+ \nu \text{ if either } \mu_+ < \nu_+ \text{ or } \mu_+ = \nu_+ \text{ and } \nu \leq \mu.$$

So the smallest element of the orbit $W\mu$ is μ_+ and if w_0 is the longest element in W , then $w_0\mu_+$ is the largest element. It is now easy to check that Dunkl operators are triangular on $\mathbb{C}[H]$ with respect to the basis e^μ partially ordered by \leq_+ . The Gram-Schmidt process applied to the monomial basis of $\mathbb{C}[H]$ yields a new basis $\{E_{\mathbf{k}}(\mu)\}_{\mu \in P}$ of $\mathbb{C}[H]$, characterized by

$$\begin{aligned} E_{\mathbf{k}}(\mu) &= e^\mu + \text{lower order terms relative to } \leq_+, \\ (E_{\mathbf{k}}(\mu), e^\nu) &= 0 \text{ for all } \nu \in P \text{ with } \nu <_+ \mu. \end{aligned}$$

Clearly, the Dunkl operators are also triangular with respect to the basis $E_{\mathbf{k}}(\mu)$ partially ordered by \leq_+ . Let $\varepsilon : \mathbb{R} \rightarrow \{\pm 1\}$ be given by $\varepsilon(x) = 1$ if $x > 0$ and $\varepsilon(x) = -1$ if $x \leq 0$. For $\mu \in P$ let $\tilde{\mu} \in \mathfrak{h}^*$ be given by

$$(2.1) \quad \tilde{\mu} = \mu + \frac{1}{2} \sum_{\alpha > 0} k_\alpha \varepsilon(\mu(\alpha^\vee)) \alpha.$$

For example for $\mu \in P_+$ dominant and regular $\tilde{\mu} = \mu + \rho_{\mathbf{k}}$, whereas for $\mu \in P_-$ antidominant $\tilde{\mu} = \mu - \rho_{\mathbf{k}}$. Then a direct computation gives for all $\mu \in P$

$$T_{\mathbf{k}}(\xi)e^\mu = \tilde{\mu}(\xi)e^\mu + \text{lower order terms},$$

which by the characterizing property of the basis $\{E_{\mathbf{k}}(\mu)\}_{\mu \in P}$ and the triangularity of $T_{\mathbf{k}}(\xi)$ yields

$$T_{\mathbf{k}}(\xi)E_{\mathbf{k}}(\mu) = \tilde{\mu}(\xi)E_{\mathbf{k}}(\mu) + \text{lower order terms}.$$

Using the hermiticity of Dunkl operators and the characterizing property of the basis $\{E_{\mathbf{k}}(\mu)\}_{\mu \in P}$ the lower order terms are easily found to vanish and we get

$$(2.2) \quad T_{\mathbf{k}}(\xi)E_{\mathbf{k}}(\mu) = \tilde{\mu}(\xi)E_{\mathbf{k}}(\mu) \text{ for all } \mu \in P.$$

Indeed triangular plus symmetric becomes diagonal. We conclude:

Theorem 2.3. *For fixed vector \mathbf{k} of coupling constants the Dunkl operators $T_{\mathbf{k}}(\xi)$ for $\xi \in \mathfrak{h}$ mutually commute. In other words, $T_{\mathbf{k}}$ extends to an algebra homomorphism $T_{\mathbf{k}} : \text{Sym}(\mathfrak{h}) \rightarrow \mathbb{D}[H^\circ] \otimes \mathbb{C}[W]$.*

We denote the coefficient of $w \in W$ in $T_{\mathbf{k}}(p)$ by $D_{\mathbf{k}}(p, w) \in \mathbb{D}(H^\circ)$, so that

$$T_{\mathbf{k}}(p) = \sum_{w \in W} D_{\mathbf{k}}(w, p) \otimes w.$$

Our choice of positive roots determines (and is determined by) a set of simple roots $\alpha_1, \dots, \alpha_n$. We shall write s_i for s_{α_i} . A rank one calculation gives for $i = 1, \dots, n$

$$s_i T_{\mathbf{k}}(\xi) - T_{\mathbf{k}}(s_i \xi) s_i = -(k_{\alpha_i} + 2k_{2\alpha_i}) \alpha_i(\xi) \text{ for all } \xi \in \mathfrak{h}$$

(recall that we put $k_\beta = 0$ if $\beta \notin R$). With induction on the degree this yields

$$s_i T_{\mathbf{k}}(p) - T_{\mathbf{k}}(s_i p) s_i = -(k_{\alpha_i} + 2k_{2\alpha_i}) T_{\mathbf{k}} \left(\frac{p - s_i p}{\alpha_i^\vee} \right) \text{ for all } p \in \text{Sym}(\mathfrak{h}).$$

In particular we find that for $q \in \text{Sym}(\mathfrak{h})^W$, $T_{\mathbf{k}}(q)$ commutes with W . So if we then write $T_{\mathbf{k}}(q) = \sum_{w \in W} D_{\mathbf{k}}(w, q) \otimes w$, then

$$(2.3) \quad D_{\mathbf{k}}(q) = \sum_{w \in W} D_{\mathbf{k}}(w, q) \in \mathbb{D}[H^\circ]^W$$

is the unique W -invariant differential operator which has the same restriction to $\mathbb{C}[H]^W$ as the Dunkl operator $T_{\mathbf{k}}(q)$. In particular, $D_{\mathbf{k}}(q)$ preserves $\mathbb{C}[H]^W$ because $T_{\mathbf{k}}(q)$ does. Since the Dunkl operators commute, we find:

Theorem 2.4. *The map $D_{\mathbf{k}} : \text{Sym}(\mathfrak{h})^W \rightarrow \mathbb{D}[H^\circ]^W$ is a \mathbb{C} -algebra homomorphism. We have $D_{\mathbf{k}}(C) = L_{\mathbf{k}} + (\rho_{\mathbf{k}}, \rho_{\mathbf{k}})$ and the image of $D_{\mathbf{k}}$ preserves $\mathbb{C}[H]^W$.*

Since the image of $D_{\mathbf{k}}$ contains the operator $L_{\mathbf{k}}$, conjugation by $\delta_{\mathbf{k}}^{\frac{1}{2}}$ of this algebra immediately gives the quantum complete integrability of the trigonometric Calogero-Moser system. The classical complete integrability follows as in the thesis of Opdam [13]. Indeed one puts $g_\alpha^2 = k_\alpha(1 - k_\alpha - 2k_{2\alpha})$ and observes that the differential operators $\delta_{\mathbf{k}}^{1/2} \circ D_{\mathbf{k}}(p) \circ \delta_{\mathbf{k}}^{-1/2}$ for $p \in \text{Sym}(\mathfrak{h})^W$ depend in a polynomial way on the g_α for $\alpha \in R$. The grading on $\text{Sym}(\mathfrak{h})$ is extended to the coupling constants $g = (g_\alpha)$ by giving each of these degree one. A symbol calculus for this total degree, which amounts to the classical limit, yields the classical complete integrability.

It was observed by Matsuo and by Cherednik that the eigenvalue problem for the Dunkl operators can be rewritten as an integrable connection of Knizhnik-Zamolodchikov type. For an exposition of this result as well as further references of the theory explained in this section we refer to the survey articles of Heckman [11] and of Opdam [24].

3. A SPECIAL EIGENSTATE OF THE TRIGONOMETRIC CALOGERO-MOSER SYSTEM

Here and in the rest of the paper we assume that R is a reduced root system. We write $k_i = k_{\alpha_i}$ for $i = 1, \dots, n$. Note that the vector $\rho_{\mathbf{k}}$ is now characterized by $\rho_{\mathbf{k}}(\alpha_i^\vee) = k_i$.

The simultaneous eigenvalue problem for the trigonometric Calogero-Moser system is obtained from the system

$$(3.1) \quad D_{\mathbf{k}}(p)F = p(\lambda)F \quad \text{for all } p \in \text{Sym}(\mathfrak{h})^W$$

by conjugation with the weight function $\delta_{\mathbf{k}}^{1/2}$. Here $\lambda \in \mathfrak{h}^*$ is fixed, and called the *spectral parameter* of the system (3.1). For the (admittedly, nonreduced) root system BC_1 , (3.1) boils down to the Euler-Gauss hypergeometric equation, and therefore (3.1) is called the *hypergeometric system* associated to the root system R . The above system has regular singular points at infinity on H° . The dimension of the local solution space equals $|W|$ and the monodromy representation of the fundamental group of $W \setminus H^\circ$, which is the affine Artin group associated with R ,

factors through the affine Hecke algebra with quadratic relations

$$(3.2) \quad (M_i - 1) \left(M_i + e^{-2\pi i k_i} \right) = 0.$$

Here M_i is the monodromy operator associated to the analytic continuation of a solution starting in the image of the positive chamber under the exponential map and making a half turn around the subtorus H_i corresponding to the simple root α_i .

Following Harish-Chandra we substitute into the hypergeometric system (3.1) a formal series of the form

$$F = \sum_{\nu \leq \mu} a_\nu e^\nu \quad \text{with} \quad a_\mu = 1.$$

With the help of (2.1) and (2.2) we see that the *exponents* $\mu \in \mathfrak{h}^*$ for which such formal solutions exist must satisfy the *indicial equation* $p(\mu + \rho_{\mathbf{k}}) = p(\lambda)$ for all $p \in S\mathfrak{h}^W$. This amounts to

$$\mu \in W\lambda - \rho_{\mathbf{k}}.$$

Choosing the base point in the positive chamber the eigenvalues of the monodromy operators of the fundamental group of H° are easily read of from these exponents. Now, if the spectral parameter $\lambda \in \mathfrak{h}^*$ is such that all points $w\lambda$ for $w \in W$ are distinct modulo the root lattice Q , i.e., if $\lambda(\alpha^\vee) \notin \mathbb{Z}$ for all $\alpha \in R$, then the hypergeometric system (3.1) has a formal solution basis. Under this nonresonance condition it is easy to see that the monodromy representation is reducible if and only if

$$(3.3) \quad \lambda(\alpha^\vee) + k_\alpha \in \mathbb{Z} \text{ for some } \alpha \in R.$$

In fact this criterion remains valid for λ resonant. We shall stick to the notation

$$(3.4) \quad \lambda = \mu + \rho_{\mathbf{k}}.$$

So μ is an exponent at infinity (in the direction of the positive chamber), whereas λ , obtained by a shift over $\rho_{\mathbf{k}}$, is the spectral parameter of our hypergeometric system.

If $\mu \in P_+$ and $k_\alpha > 0$ for all $\alpha \in R$, then the hypergeometric system (3.1) has solutions in $\mathbb{C}[H]^W$ with leading exponent $\mu = \lambda - \rho_{\mathbf{k}}$. This is an extreme case of reducibility since (remember that $\rho_{\mathbf{k}}(\alpha_i^\vee) = k_i$) we have $\lambda(-\alpha_i^\vee) + k_i = \mu(-\alpha_i^\vee) \in -\mathbb{N}$ for $i = 1, \dots, n$. Hence for generic \mathbf{k} the exponents $\mu \in P_+$ are intersection points of normal crossing reducibility hyperplanes (3.3). Clearly the solutions in $\mathbb{C}[H]^W$ are fixed under the monodromy representation.

The goal of this section is to discuss the next best case for which the hypergeometric system admits a monodromy invariant subspace of solutions of dimension $n + 1$ that affords the reflection representation of the affine Artin group as its monodromy. These special solutions of (3.1) satisfy a bigger set of differential equations, a set which we baptize the *special hypergeometric system*. It turns out to be generated by quadratic equations and to give rise to a projective structure on H° .

From now on we assume that the reduced root system R is irreducible. We denote by $\varpi_1, \dots, \varpi_n$ the fundamental weights in P_+ . In dealing with particular cases the numbering will be the same as in the Bourbaki text on root systems [5]. This means in particular that $R = W\{\alpha_1, \alpha_n\}$.

We shall describe for each root system a sequence $(\mu_1, \dots, \mu_{n+1})$ in \mathfrak{h}^* given up to order (repetition is not excluded) which describes the exponents of the special hypergeometric system with multiplicities. This will be done in a case by case manner. The points $\mu_i = \mu_{i,\mathbf{k}}$ will depend linearly on \mathbf{k} and in the case of type A_n on an additional parameter as well. Since \mathbf{k} takes on R at most two values, namely k_1 and k_n , we find it convenient to adapt our notation accordingly: we write k for k_1 and if $\alpha_n \notin W\alpha_1$ (i.e., R of type B, C, F or G), k' stands for k_n . In the A_n -case, k' will have a different meaning, but we will then abuse notation, by letting \mathbf{k} then stand for the pair $(k = k_1, k')$. We shall first discuss a root system whose Dynkin diagram is a chain, i.e., one of type A, B, C, F or G , and subsequently deal with those of type D or E .

For R of type A_n ($n \geq 2$), the extended Dynkin diagram has a loop causing an additional modulus in the affine reflection representations. This accounts for an additional parameter $k' \in \mathbb{C}$. We put

$$(3.5) \quad \mu_i = (x - ik)\varpi_{i-1} + ((i-1)k - x)\varpi_i$$

for $i = 1, \dots, n+1$ with the convention $\varpi_0 = \varpi_{n+1} = 0$. Here we take $x, y \in \mathbb{C}$ related to $k, k' \in \mathbb{C}$ by

$$x + y = (n+1)k, \quad x - y = (n+1)k'$$

so that $\mu_1 = -x\varpi_1, \mu_{n+1} = -y\varpi_n$. Since $\alpha_i = -\varpi_{i-1} + 2\varpi_i - \varpi_{i+1}$ it is clear that

$$\mu_{i+1} - \mu_i = (x - ik)\alpha_i \quad \text{for } i = 1, \dots, n.$$

Keeping in mind (3.4) it is clear from (3.5) that $\lambda_i(\alpha_i^\vee) = ik - x$ and so the last equality yields

$$\lambda_{i+1} = s_i(\lambda_i) \quad \text{for } i = 1, \dots, n.$$

For R of type B_n we still take (3.5) for $i = 1, \dots, n-1$ and in accordance with (3.4) and the above identity, we put

$$\begin{aligned} \mu_n &= \mu_{n-1} - ((n-1)k - x)\alpha_{n-1} \\ &= (x - nk)\varpi_{n-1} + (2(n-1)k - 2x)\varpi_n, \\ \mu_{n+1} &= \mu_n - (2(n-1)k + k' - 2x)\alpha_n \\ &= ((n-2)k + k' - x)\varpi_{n-1} + (2x - 2(n-1)k - 2k')\varpi_n, \end{aligned}$$

with the convention (for R of type B, C, F or G) that $k' = k_n$. Hence if we take $x = (n-2)k + k', y = 2k$ then we find $\mu_1 = -x\varpi_1$ and $\mu_{n+1} = -y\varpi_n$.

Likewise for R of type C_n we find

$$\begin{aligned} \mu_n &= \mu_{n-1} - ((n-1)k - x)\alpha_{n-1} \\ &= (x - nk)\varpi_{n-1} + ((n-1)k - x)\varpi_n, \\ \mu_{n+1} &= \mu_n - ((n-1)k + k' - x)\alpha_n \\ &= ((n-2)k + 2k' - x)\varpi_{n-1} + (x - (n-1)k - 2k')\varpi_n, \end{aligned}$$

and with $x = (n-2)k + 2k', y = k$ we get $\mu_1 = -x\varpi_1$ and $\mu_{n+1} = -y\varpi_n$.

In case R is of type F_4 we take

$$\begin{aligned}\mu_1 &= -2x\alpha_1 - 3x\alpha_2 - 4x\alpha_3 - 2x\alpha_4 = -(k + k')\varpi_1, \\ \mu_2 &= -y\alpha_1 - 3x\alpha_2 - 4x\alpha_3 - 2x\alpha_4 = (-k + k')\varpi_1 - k'\varpi_2, \\ \mu_3 &= -y\alpha_1 - 2y\alpha_2 - 4x\alpha_3 - 2x\alpha_4 = (-2k + k')\varpi_2 + (2k - 2k')\varpi_3, \\ \mu_4 &= -y\alpha_1 - 2y\alpha_2 - 3y\alpha_3 - 2x\alpha_4 = -2k\varpi_3 + (2k - k')\varpi_4, \\ \mu_5 &= -y\alpha_1 - 2y\alpha_2 - 3y\alpha_3 - y\alpha_4 = -(2k + k')\varpi_4,\end{aligned}$$

with $x = k + k'$, $y = 2k + k'$. Note that $\mu_2 - \mu_1 = k'\alpha_1$, $\mu_3 - \mu_2 = (-k + k')\alpha_2$, $\mu_4 - \mu_3 = (-2k + k')\alpha_3$ and $\mu_5 - \mu_4 = -2k\alpha_4$.

Finally for R of type G_2 , let

$$\begin{aligned}\mu_1 &= -2x\alpha_1 - x\alpha_2 = -x\varpi_1, \\ \mu_2 &= -3y\alpha_1 - x\alpha_2 = (x - 2k)\varpi_1 + (k - x)\varpi_2, \\ \mu_3 &= -3y\alpha_1 - 2y\alpha_2 = -y\varpi_2,\end{aligned}$$

with $x = \frac{1}{2}(k + 3k')$, $y = \frac{1}{2}(k + k')$. Note that $\mu_2 - \mu_1 = \frac{1}{2}(-k + k')\alpha_2$. This completes our case by case descriptions of the set of special exponents for R of type A, B, C, F or G .

For R of type D_n we take for $i = 1, \dots, n - 2$

$$\mu_i = (n - 2 - i)k\varpi_{i-1} - (n - 1 - i)k\varpi_i, \quad \mu_{n-1} = -2k\varpi_{n-1}, \quad \mu_n = -2k\varpi_n,$$

where $\varpi_0 = 0$, as before. With $x = (n - 2)k$ and $y = 2k$ we have $\mu_1 = -x\varpi_1$, $\mu_{n-1} = -y\varpi_{n-1}$, $\mu_n = -y\varpi_n$.

For R of type E_8 let

$$\begin{aligned}\mu_1 &= -3k\varpi_1, & \mu_5 &= -2k\varpi_5 + k\varpi_6, \\ \mu_2 &= -2k\varpi_2, & \mu_6 &= -3k\varpi_6 + 2k\varpi_7, \\ \mu_3 &= k\varpi_1 - 2k\varpi_3, & \mu_7 &= -4k\varpi_7 + 3k\varpi_8, \\ \mu_4 &= -k\varpi_4, & \mu_8 &= -5k\varpi_8.\end{aligned}$$

The case R of type E_7 is obtained by forgetting μ_8 and putting $\varpi_8 = 0$ in the expression for μ_7 and likewise for E_6 . Note that for both D_n and E_n

$$\lambda(\alpha_i^\vee) = -d(i)k$$

with $d(i)$ the distance in the Dynkin diagram from the i^{th} node to the triple node. For type D_n and E_n put $\mu_{n+1} = \mu_{n-2}$ and $\mu_{n+1} = \mu_4$ respectively, so μ_i occurs with multiplicity two if i labels the triple node. All in all, we have now given for each irreducible root system R of rank n a set $\{\mu_1, \dots, \mu_{n+1}\}$ in \mathfrak{h}^* depending linearly on k resp. on (k, k') depending on whether or not the type is D or E . We shall write K for the space of coupling parameters, so $k \in K = \mathbb{C}$ or $(k, k') \in K = \mathbb{C}^2$ respectively.

Definition 3.1. *The special exponents of R are the terms of $(\mu_1, \dots, \mu_{n+1})$ in \mathfrak{h}^* , counted with multiplicity.*

Notice that each μ_i depends linearly on $\mathbf{k} \in K$. We denote the inner product on a dual to $(,)$ by $(,)^\vee$ and write C^\vee for the corresponding element of $\text{Sym}^2(\mathfrak{h}^*)$.

For the A_n root system we define for every positive root $\alpha \in \mathfrak{a}^*$ an element $\alpha' \in \mathfrak{a}$ as the vector in the characterized by the property that $\beta(\alpha') \geq 0$ for all $\beta \in R_+$ with equality when $\beta = \alpha$ or $(\beta, \alpha) = 0$ and with $(\alpha', \alpha)^\vee(\alpha, \alpha) = 4\frac{n-1}{n+1}$. In terms of the standard description in Bourbaki [5]: if $\alpha^\vee = e_i - e_j$ with $i < j$, then α' is the image of $e_i + e_j$ in $\mathbb{R}^{n+1}/\text{main diagonal}$.

Proposition 3.2. *The special exponents are exactly the solutions in μ of the following quadratic equation in $\text{Sym}^2(\mathfrak{h}^*)$:*

$$\mu^2 + \frac{1}{2} \sum_{\alpha > 0} \mu(k\alpha^\vee + k'\alpha')\alpha^2 + a(k, k')C^\vee = 0 \quad \text{for type } A_n, n \geq 2,$$

$$\mu^2 + \frac{1}{2} \sum_{\alpha > 0} \mu(k_\alpha\alpha^\vee)\alpha^2 + a(k, k')C^\vee = 0 \quad \text{for type } B, C, F \text{ or } G,$$

$$\mu^2 + \frac{1}{2} \sum_{\alpha > 0} \mu(k_\alpha\alpha^\vee)\alpha^2 + a(k)C^\vee = 0 \quad \text{for type } D \text{ or } E.$$

The function a is the quadratic polynomial in the coupling constants: for R of type A, B, C, F or G it is given by $a(k, k') = xy(\varpi_1, \varpi_n) = (\mu_1, \mu_{n+1})$ and for type D_n, E_6, E_7, E_8 (with inner product normalized by $(\alpha, \alpha) = 2$) by $a(k) = (n-2)k^2$ and $a(k) = 6k^2, 12k^2, 30k^2$ respectively.

For $\mathbf{k} \in K$, we define a linear map

$$(3.6) \quad D_{\mathbf{k}} : \text{Sym}^2(\mathfrak{h}) \rightarrow \mathbb{D}[H^\circ]$$

by letting its value on the monomial $\xi\eta$ (with $\xi, \eta \in \mathfrak{h}$) be

$$\partial(\xi)\partial(\eta) + \frac{1}{2} \sum_{\alpha > 0} \alpha(\xi)\alpha(\eta) \frac{1 + e^{-\alpha}}{1 - e^{-\alpha}} k_\alpha \partial(\alpha^\vee) + \rho_{\mathbf{k}}(\xi)\rho_{\mathbf{k}}(\eta)$$

if the root system is not of type A_n and

$$\partial(\xi)\partial(\eta) + \frac{1}{2} \sum_{\alpha > 0} \alpha(\xi)\alpha(\eta) \left(\frac{1 + e^{-\alpha}}{1 - e^{-\alpha}} k \partial(\alpha^\vee) + k' \partial(\alpha') \right) + \rho_{\mathbf{k}}(\xi)\rho_{\mathbf{k}}(\eta)$$

if the root system is of type $A_n, n \geq 2$. Observe that the notations in (2.3) and (3.6) are compatible for $p \in \text{Sym}^2(\mathfrak{h})^W$.

Theorem 3.3. *The system of differential equations*

$$(3.7) \quad D_{\mathbf{k}}(p)F = (p(\rho_{\mathbf{k}}) - a(\mathbf{k})(C^\vee, p))F \quad \text{for all } p \in \text{Sym}^2(\mathfrak{h}),$$

enjoys the following properties:

- (i) *the system is integrable on H° and invariant under W ,*
- (ii) *solutions of (3.7) are also solutions of the general hypergeometric system (3.1) with spectral parameter $\lambda_i = \mu_i(\mathbf{k}) + \rho_{\mathbf{k}}$, where μ_i is a special exponent,*

- (iii) the dimension of the local solution space of (3.7) is equal to $n + 1$, and the monodromy of the associated rank $n + 1$ local system over $W \setminus H^\circ$ factors through the reflection representation of the affine Hecke algebra (viewed as a quotient of the group algebra of the fundamental group of $W \setminus H^\circ$): M_i has eigenvalue 1 with multiplicity n and eigenvalue $-e^{2\pi i k_i}$ with multiplicity 1, in accordance with (3.2).

We call (3.7) the *special hypergeometric system* with parameter \mathbf{k} .

Theorem 3.4. *For real $\mathbf{k} \in K$ the special hypergeometric system (3.7) admits a monodromy invariant hermitian form $h_{\mathbf{k}}$. This is nondegenerate and has Lorentzian signature if $\mathbf{k} \in K_+$, where*

$$K_+ := \begin{cases} (0, \frac{1}{n-2}) & \text{for type } D_n, \\ (0, \frac{1}{n-3}) & \text{for type } E_n, \\ \{(k, k') \in (-\frac{1}{2}, \frac{1}{2})^2 : 0 < x < 1; 0 < y < 1\} & \text{otherwise.} \end{cases}$$

Lemma 3.5. *For a basis F_0, \dots, F_n of a local solution space of (3.7) and a basis ξ_1, \dots, ξ_n of \mathfrak{h} the Wronskian is (up to a constant depending on choices of bases) given by*

$$\det \begin{pmatrix} F_0 & F_1 & \dots & F_n \\ \partial(\xi_1)F_0 & \partial(\xi_1)F_1 & \dots & \partial(\xi_1)F_n \\ \vdots & \vdots & \dots & \vdots \\ \partial(\xi_n)F_0 & \partial(\xi_n)F_1 & \dots & \partial(\xi_n)F_n \end{pmatrix} = \delta_{\mathbf{k}}^{-1}.$$

Fix a base point $z_0 \in W \setminus H^\circ$, and let $\mathbf{k} \in K_+$. We denote the projectivized monodromy group of (3.7) by $\Gamma = \Gamma_{\mathbf{k}}$ and by $\widetilde{W \setminus H^\circ} \rightarrow W \setminus H^\circ$ the corresponding Γ -covering. It follows from Lemma 3.5 that the projectivized evaluation map given by $[F_0(z) : \dots : F_n(z)]$ becomes a local isomorphism

$$\text{Pev} : \widetilde{W \setminus H^\circ} \longrightarrow \mathbb{P}^n(\mathbb{C})$$

and so equation (3.7) gives a projective structure on $W \setminus H^\circ$. The monodromy invariant hermitian form yields a complex hyperbolic ball $\mathbb{B}^n \subset \mathbb{P}^n$ on which Γ acts by biholomorphic automorphisms.

Definition 3.6. *We say that $\mathbf{k} \in K_+$ is admissible if the image of Pev is contained in \mathbb{B}^n , so that we have a map $\text{Pev} : \widetilde{W \setminus H^\circ} \rightarrow \mathbb{B}^n$.*

Admissibility was obtained for the additive setting [7] and we expect that its proof can be carried over to the multiplicative setting. For the rest of this section we assume that the parameter \mathbf{k} is admissible.

It is our goal to find conditions on \mathbf{k} such that Γ acts on \mathbb{B}^n as a lattice group in $PU(n, 1)$, so that the commutative diagram

$$(3.8) \quad \begin{array}{ccc} \widetilde{W \setminus H^\circ} & \xrightarrow{\text{Pev}} & \mathbb{B}^n \\ \downarrow & & \downarrow \\ W \setminus H^\circ & \longrightarrow & \Gamma \backslash \mathbb{B}^n \end{array}$$

lives in the complex-analytic category. Notice that in the bottom line the left hand side is in the realm of *projective* geometry, whereas the right hand side concerns *hyperbolic* geometry. Even if Γ does not act as a lattice group on \mathbb{B}^n the space $W \setminus H^\circ$ always acquires a hyperbolic metric under the admissibility assumption $\mathbf{k} \in K_+$. If it happens to be complete, then $\text{Pev} : \widetilde{W \setminus H^\circ} \rightarrow \mathbb{B}^n$ is a covering map, and hence a holomorphic isomorphism, since \mathbb{B}^n is simply connected. So $W \setminus H^\circ$ has then the structure of a ball quotient.

However it is rare that the hyperbolic metric on $W \setminus H^\circ$ is complete. In the one-dimensional case of the Euler-Gauss hypergeometric equation there is only one case, namely when all exponent differences are zero and Pev gives the uniformization of $\mathbb{P}^1 - \{0, 1, \infty\}$ by the unit disc. To remedy the above argument one seeks a (partial) stratified compactification on the projective side of diagram (3.8) and an extension of Pev over the associated boundary as a locally biholomorphic map, such that the hyperbolic metric on the compactified space is complete. Now one can argue as before. To find the correct compactification on the projective side is a very delicate problem in birational geometry, and its mere existence imposes strong restrictions on $\mathbf{k} \in K_+$, called the *Schwarz conditions* in the additive case [7]. A direct approach to find the desired compactification is due to Deligne and Mostow [8], and uses the geometric invariant theory of binary forms. Its scope is limited to the case of classical root systems, but it does construct on the projective geometric side the compactification that corresponds to the Baily-Borel compactification on the ball quotient side.

An alternative approach, followed in [7], is to work with a rather easily described compactification that dominates the searched compactification. The map Pev extends to a holomorphic submersion onto \mathbb{B}^n with connected components of fibers being compact, and one can apply Stein factorization to it; the Stein factor then gives the desired compactification. Although this procedure is hardly constructive in the projective category, the outcome is that the complement of the image of Pev in \mathbb{B}^n is a locally finite union of hyperballs (a so-called Heegner divisor), and typically the group Γ contains reflections with the hyperballs as mirrors.

Question 3.7. Is there a modular interpretation of diagram (3.8), such that the projectivized evaluation map Pev gets the interpretation of a period map?

For R a classical root system the answer to the question is positive, and given by the theory of Deligne and Mostow [8]. For R of type E one encounters moduli spaces of Del Pezzo surfaces and suitable period maps on these [1], [15], [12]. But the picture for general root systems R remains unclear. In the next section we shall discuss the case R of type E_8 with $k = \frac{1}{6}$.

4. AN EXAMPLE: THE E_8 ROOT SYSTEM WITH COUPLING CONSTANT $\frac{1}{6}$

The Schwarz conditions for the classical Euler-Gauss hypergeometric equation require that at the three singular points $0, 1, \infty$ the exponent differences are the reciprocal of integers. Once the correct stratified compactification of $W \setminus H^\circ$ adapted

to the multivariable special hypergeometric system (3.7) is found the Schwarz conditions boil down to similar conditions as for the one dimensional case, transverse to the codimension strata of the boundary. The special hypergeometric system (3.7) was given as a system on H° invariant under W . However from Theorem (3.3) it is clear that the system lives in fact on the adjoint torus H' (with character lattice $Q = \mathbb{Z}R$), and is invariant under $W' = W \rtimes \text{Aut}(R_+)$ for R not of type A_n . For type A_n we have to impose the condition $k' = 0 \iff x = y$ so $-w_0\mu_1 = \mu_{n+1}$. In principle working on the quotient space $W' \backslash H'^\circ$ of $W \backslash H^\circ$ allows more solutions to the Schwarz conditions. However in certain examples there is still more symmetry which is unclear from the root system perspective. One ought to think of this as a *hidden symmetry* for the system (3.7).

Example 4.1. For the root system R of type A_n one has a covering map

$$W' \backslash H'^\circ \longrightarrow \mathcal{D}_{n+3}$$

of degree $\binom{n+3}{2}$ with \mathcal{D}_{n+3} the moduli space of reduced effective degree $(n+3)$ divisors on the projective line \mathbb{P} . Indeed $H^\circ = \{(x_1, \dots, x_{n+1}) \in (\mathbb{C}^\times)^{n+1}; x_1 \dots x_{n+1} = 1, x_i \neq x_j \text{ for } 1 \leq i < j \leq n+1\}$ and $H'^\circ = C_{n+1} \backslash H^\circ$ with C_{n+1} the group of roots of unity of order $(n+1)$ acting diagonally, while $W = S_{n+1}$ and $W' = S_{n+1} \rtimes S_2$. Writing $H'^\circ = \mathbb{C}^\times \backslash \{(x_0 = 0, x_1, \dots, x_{n+1}, x_{n+2} = \infty) \in \mathbb{P}^{n+3}; x_i \neq x_j \text{ for } 0 \leq i < j \leq n+2\}$ the map $(x_0, \dots, x_{n+2}) \mapsto \{x_0, \dots, x_{n+2}\}$ yields the above covering map. The special hypergeometric system is the pull back of the Lauricella system on \mathcal{D}_{n+3} if $(n+3)k = 2$. Under this hypothesis, the signature of the invariant hermitian form is Lorentzian. The Schwarz conditions amount to $(\frac{1}{2} - k)^{-1} = q \in \mathbb{N}$. This leaves us with the following possibilities

n	1	2	3	5	9
q	∞	10	6	4	3

which are exactly those cases in the table of Deligne and Mostow with all Lauricella exponents equal (to k). In all cases with $q \in 2\mathbb{N}$ the Schwarz condition would already be satisfied on $\mathcal{M}_{0,n+3}$. The quotient by S_{n+3} is therefore only needed to include the last case $n = 9, q = 3$ which gives

$$(4.1) \quad \mathcal{D}_{12} \hookrightarrow \Gamma_9 \backslash \mathbb{B}^9$$

with Γ_9 generated by order 3 hyperbolic reflections. Their mirrors are hyperballs in \mathbb{B}^9 whose quotient by Γ_9 is exactly the boundary of \mathcal{D}_{12} under the map (4.1). For $D \in \mathcal{D}_{12}$ represented by $\{x_0, \dots, x_{11}\}$ consider the algebraic curve C with affine equation $y^6 = \prod_{i=0}^{11} (x - x_i)$. The periods

$$\int_{x_i}^{x_{i+1}} \frac{dx}{y}$$

are Lauricella integrals and solutions of (3.7) for R of type A_9 with $(k, k') = (\frac{1}{6}, 0)$.

A rational elliptic surface S in Weierstrass form is a projective surface with affine equation

$$(4.2) \quad S : y^2 = 4x^3 - g_2x - g_3$$

with $g_2, g_3 \in \mathbb{C}[u, v]$ binary forms of degree 4 and 6 respectively. The surface represents an elliptic fibration over the projective line \mathbb{P} with coordinate $t = u/v$. The *discriminant divisor* D of the fibration is the divisor on that line defined by $\Delta := g_2^3 - 27g_3^2 \in \mathbb{C}[u, v]$; its support gives the singular fibers of the fibration, but D gives each such fiber a multiplicity. Since Δ is a binary form of degree 12, D is an effective degree 12 divisor on \mathbb{P} . We have a zero of multiplicity 1 precisely when the corresponding singular fiber is a nodal rational curve. Let \mathcal{M} be the moduli space of rational elliptic surfaces of the form (4.2) with reduced discriminant divisor. There is a natural map

$$\mathcal{M} \rightarrow \mathcal{D}_{12}, \quad S \mapsto D$$

which turns out to be injective. Since $\dim \mathcal{M} = 8$ and $\dim \mathcal{D}_{12} = 9$ the space \mathcal{M} lies as a hypersurface in \mathcal{D}_{12} . The main result of [12] is the following theorem.

Theorem 4.2. *We have a commutative diagram*

$$\begin{array}{ccc} \mathcal{M} & \longrightarrow & \Gamma_8 \backslash \mathbb{B}^8 \\ \downarrow & & \downarrow \\ \mathcal{D}_{12} & \longrightarrow & \Gamma_9 \backslash \mathbb{B}^9 \end{array}$$

with the bottom arrow the Deligne-Mostow period map (4.1).

Relative to the natural line bundle on \mathcal{D}_{12} , the hypersurface \mathcal{M} in \mathcal{D}_{12} has degree 3762 [25], and so from a projective point of view \mathcal{M} lies in a wildly bended manner in \mathcal{D}_{12} . On the other hand, through hyperbolic eyes the embedding is totally geodesic (which can be thought of as linear). To understand the connection of Theorem 4.2 and the previous section we recall a result of Looijenga [16].

Theorem 4.3. *Let \mathcal{M}' be the moduli space of rational elliptic surfaces of the form (4.2) with reduced discriminant, marked with a distinguished singular nodal rational fiber. Then we have a commutative diagram*

$$(4.3) \quad \begin{array}{ccc} \mathcal{M}' & \xrightarrow{\cong} & W \backslash H^\circ \\ \searrow & & \swarrow \\ & \mathcal{M} & \end{array}$$

with $\mathcal{M}' \rightarrow \mathcal{M}$ the natural degree 12 covering map given by forgetting the distinguished fiber, and $W \backslash H^\circ$ the regular toric orbit space for R of type E_8 .

The description of the isomorphism (4.3) is in geometric terms. There is a distinguished section in the Weierstrass model S over \mathbb{P} through $(x : y : z) = (0 : 1 : 0)$ called the zero section. The global sections of S over \mathbb{P} form a lattice by addition in fibers, and with respect to minus the intersection form on S , is called the *Mordell-Weil lattice*. It turns out for S in \mathcal{M} that the Mordell-Weil lattice is positive definite, even, unimodular and of rank 8, and hence isomorphic to the E_8 -lattice. From the group law on the regular fibers the distinguished singular nodal rational fiber minus

the node inherits a group law isomorphic to the multiplicative group \mathbb{C}^\times . Looking at intersection points of the Mordell-Weil lattice with the distinguished singular fiber gives the isomorphism (4.3).

Combining the above two theorems gives a geometric interpretation of diagram (3.8) for R of type E_8 and $k = \frac{1}{6}$.

Theorem 4.4. *We have a commutative diagram*

$$\begin{array}{ccc} \mathcal{M}' & \longrightarrow & \Gamma'_8 \backslash \mathbb{B}^8 \\ \downarrow & & \downarrow \\ \mathcal{M} & \longrightarrow & \Gamma_8 \backslash \mathbb{B}^8 \end{array}$$

with the bottom arrow given by the top arrow of the diagram in Theorem (4.2) and the top arrow being the bottom arrow of diagram (3.8) with R of type E_8 and $k = \frac{1}{6}$ in accordance with diagram (4.3).

Remark 4.5. For R of type E_8 with Coxeter number $h = 30$, the exponent difference along the blown up identity element equals $1 - hk = 1 - 30/6 = -4$. Since $w_0 = -1$ the exponent difference on the quotient by W of this divisor becomes -2 which is not the reciprocal of an integer. Apparently the Schwarz conditions only hold after exploitation of the hidden symmetry on $W \backslash H^\circ$. Since Γ'_8 need not be a normal subgroup of Γ_8 this hidden symmetry is just an equivalence relation on $W \backslash H^\circ$ with classes of 12 element sets, and need not be obtained from an action of a group of order 12 on $W \backslash H^\circ$. A description of this hidden symmetry in terms of the E_8 root system remains unclear.

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