

A discrete entropic uncertainty relation

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Recently [MaU] a new class of ‘generalised entropic’ uncertainty relations for the probability distributions of non-commuting random variables was proved as a simple consequence of the Riesz-Thorin interpolation theorem. Here we shall give a quite explicit proof of the central inequality of this class, an ‘entropic’ uncertainty relation, which has been conjectured by Kraus [Kra].

We consider the following situation, not uncommon in quantum mechanics. Two observables of a physical system are represented by symmetric complex $n \times n$ matrices A and B , which we shall assume to have non-degenerate spectra. We can write A and B in the form

$$A = \sum_{i=1}^n \alpha_i P_i \quad \text{and} \quad B = \sum_{i=1}^n \beta_i Q_i,$$

where P_1, \dots, P_n and Q_1, \dots, Q_n are sequences of mutually orthogonal one-dimensional projections, and the sequences $\alpha_1, \dots, \alpha_n$ and β_1, \dots, β_n consist of distinct real numbers, to be interpreted as the values which the observables can take. Each state ω on the algebra M_n of all complex $n \times n$ matrices then induces probability distributions on the spectra of A and B : $\omega(P_i)$ (or $\omega(Q_i)$) is the probability to find the value α_i (or β_i) when measuring the observable A (or B). One now defines the *uncertainty* $H(A, \omega)$ of A in the state ω as the Shannon entropy of this probability distribution:

$$H(A, \omega) = - \sum_{i=1}^n \omega(P_i) \log \omega(P_i).$$

The question was raised ([BBM], [Deu], [Kra]), what can be said about $H(A, \omega) + H(B, \omega)$, more in particular about its lower bound

$$d(A, B) = \inf_{\omega} (H(A, \omega) + H(B, \omega)).$$

*Supported by the Netherlands Organisation of Scientific Research NWO.

One may regard this infimum as a "degree of incompatibility" of the observables A and B .

As a first reduction, let us note that $H(A, \omega)$ does not depend on the real numbers $\alpha_1, \dots, \alpha_n$, but only on the projections P_1, \dots, P_n , which are the minimal projections in the maximal abelian von Neumann algebra \mathcal{A} generated by A . Let us therefore write $H(\mathcal{A}, \omega)$, $H(\mathcal{B}, \omega)$ and $d(\mathcal{A}, \mathcal{B})$ in what follows. When viewed in this way, d becomes a natural distance function between maximal abelian von Neumann algebras, comparable to the distance of point *sets* (not of points) in geometry: $d(\mathcal{A}, \mathcal{B}) = 0$ if and only if \mathcal{A} and \mathcal{B} have a minimal projection in common.

The latter observation suggests to consider the following easily computable functional on pairs of abelian von Neumann algebras in M_n :

$$m(\mathcal{A}, \mathcal{B}) = \max\{\text{tr}PQ \mid P \in \mathcal{A}, Q \in \mathcal{B} \text{ minimal projections}\}.$$

This definition amounts to

$$m(\mathcal{A}, \mathcal{B}) = \max_{1 \leq i, j \leq n} \text{tr}(P_i Q_j) = \max_{1 \leq i, j \leq n} |\langle e_i, f_j \rangle|^2,$$

where e_i and f_j are unit vectors in the ranges of P_i and Q_j respectively. Note that

$$m(\mathcal{A}, \mathcal{B}) \leq 1,$$

with equality if and only if $d(\mathcal{A}, \mathcal{B}) = 0$. On the other hand, since for all j

$$\sum_{i=1}^n \text{tr}(P_i Q_j) = \text{tr}Q_j = 1,$$

we have

$$m(\mathcal{A}, \mathcal{B}) \geq \frac{1}{n}.$$

It was observed by Kraus [Kra] that this lower bound is reached for 'complementary' observables A and B , which corresponds to e_j and f_j of the form

$$(e_j)_k = 1 \text{ if } j = k, 0 \text{ otherwise;}$$

$$(f_j)_k = \frac{1}{\sqrt{n}} e^{\frac{2\pi i j k}{n}}.$$

Note that the algebras \mathcal{A} and \mathcal{B} then take the form:

$$\mathcal{A} = \{X \in M_n \mid X \text{ diagonal}\},$$

$$\mathcal{B} = \{X \in M_n \mid X_{i+k, j+k} = X_{i, j} \text{ for all } i, j, k\},$$

where the addition of indices is taken modulo n .

Kraus went on to conjecture that for such complementary observables (or algebras in our terminology)

$$d(\mathcal{A}, \mathcal{B}) = \log n.$$

Indeed, the inequality $d(\mathcal{A}, \mathcal{B}) \leq \log n$ is easily established: choose $\omega(X) = \langle e_1, X e_1 \rangle$, so that $(\omega(P_1), \dots, \omega(P_n)) = (1, 0, \dots, 0)$ and $(\omega(Q_1), \dots, \omega(Q_n)) = (\frac{1}{n}, \frac{1}{n}, \dots, \frac{1}{n})$; then $H(\mathcal{A}, \omega) = 0$ and

$$H(\mathcal{B}, \omega) = - \sum_{i=1}^n \frac{1}{n} \log \frac{1}{n} = \log n.$$

Kraus' conjecture is therefore a consequence of the following theorem.

Theorem 1 *For all maximal abelian von Neumann subalgebras \mathcal{A} and \mathcal{B} of M_n one has*

$$d(\mathcal{A}, \mathcal{B}) \geq -\log m(\mathcal{A}, \mathcal{B}).$$

Proof. Let $\{P_1, \dots, P_n\}$ and $\{Q_1, \dots, Q_n\}$ be complete sets of minimal projections in \mathcal{A} and \mathcal{B} respectively, and let e_j and f_j be unit vectors in the ranges of P_j and Q_j respectively ($j = 1, \dots, n$). We may assume that $\{e_j\}$ is the canonical basis of \mathbf{C}^n . From the concavity of the function $\eta : [0, 1] \rightarrow [0, \infty) : x \mapsto -x \log x$ (with $\eta(0) := 0$) it follows that the minimum of $H(\mathcal{A}, \omega) + H(\mathcal{B}, \omega) = \sum_{j=1}^n (\eta(\omega(P_j)) + \eta(\omega(Q_j)))$ is taken in a vector state $\omega(X) = \langle \psi, X \psi \rangle$ on M_n . It therefore suffices to prove that for all $\psi \in \mathbf{C}^n$:

$$\sum_{j=1}^n (\eta(|\langle e_j, \psi \rangle|^2) + \eta(|\langle f_j, \psi \rangle|^2)) \geq -\log \max_{i,j} |\langle e_i, f_j \rangle|^2. \quad (1)$$

Now let $m = m(\mathcal{A}, \mathcal{B}) = \max_{i,j} |\langle e_i, f_j \rangle|^2$ and let a unitary map $T : \mathbf{C}^n \rightarrow \mathbf{C}^n$ be defined by $T f_j = e_j$. Then $\psi_i = \langle e_i, \psi \rangle$ and $(T\psi)_i = \langle e_i, T\psi \rangle = \langle T^{-1} e_i, \psi \rangle = \langle f_i, \psi \rangle$. If we now write $h(\psi)$ for $\sum_{j=1}^n \eta(|\psi_j|^2)$, then the inequality takes the form

$$h(\psi) + h(T\psi) \geq -\log m. \quad (2)$$

For $n \in \mathbf{N}$ and $p \in [1, \infty]$ let $l^p(n)$ denote the Banach space \mathbf{C}^n with norm

$$\|\psi\|_p = \begin{cases} (\sum_{i=1}^n |\psi_i|^p)^{\frac{1}{p}} & \text{if } 1 \leq p < \infty \\ \max_{1 \leq i \leq n} |\psi_i| & \text{if } p = \infty. \end{cases}$$

We now make the following observation.

Lemma 2 *For all unit vectors ψ in \mathbf{C}^n*

$$\frac{d}{dp^{-1}} \|\psi\|_p|_{p=2} = h(\psi).$$

Proof. First we note that

$$\begin{aligned} \frac{d}{dp} \|\psi\|_p^p|_{p=2} &= \frac{d}{dp} \sum_{j=1}^n |\psi_j|^p|_{p=2} = \\ \sum_{j=1}^n |\psi_j|^2 \log |\psi_j| &= -\frac{1}{2} \sum_{j=1}^n \eta(|\psi_j|^2) = -\frac{1}{2} h(\psi). \end{aligned}$$

Therefore, since $\|\psi\|_2 = 1$,

$$\begin{aligned} \frac{d}{dp^{-1}} \|\psi\|_p|_{p=2} &= \frac{d}{dp^{-1}} (\|\psi\|_p^p)^{p^{-1}}|_{p=2} = \\ \log \|\psi\|_2^2 + \frac{1}{p} (\|\psi\|_p^p)^{\frac{1}{p}-1} \cdot \frac{dp}{dp^{-1}} \cdot \frac{d}{dp} \|\psi\|_p^p|_{p=2} & \\ = \frac{1}{2} \cdot (-4) \left(-\frac{1}{2} h(\psi)\right) &= h(\psi). \end{aligned}$$

□

Let $\|T\|_p$ denote the norm of the linear map $T : \mathbf{C}^n \rightarrow \mathbf{C}^n$, viewed as an operator $l^p(n) \rightarrow l^q(n)$, where $\frac{1}{p} + \frac{1}{q} = 1$. (Here we make the usual convention that $\frac{1}{\infty} = 0$.) Then $\|T\|_2 = 1$ and $\|T\|_1 = \max_{j,k} |T_{jk}| = \max_{j,k} |\langle f_j, e_k \rangle| = \sqrt{m}$.

Theorem 3 (Riesz-Thorin interpolation) *For a linear map $T : \mathbf{C}^n \rightarrow \mathbf{C}^n$ the function*

$$f_T : [0, 1] \rightarrow \mathbf{R} : \frac{1}{p} \mapsto \log \|T\|_p$$

is convex.

Proof: [Rie]; see also [HLP]. □

It follows that f_T has a right derivative $f'_T(\frac{1}{2})$ at $\frac{1}{2}$, and that, since $f_T(\frac{1}{2}) = \log \|T\|_2 = 0$ and $f_T(1) = \log \|T\|_1 = \frac{1}{2} \log m$, we have

$$f'_T(\frac{1}{2}) \leq \frac{f_T(1) - f_T(\frac{1}{2})}{1 - \frac{1}{2}} = \log m.$$

On the other hand, by the definition of the operator norm $\|T\|_p$, we have for all $p \in [1, \infty]$ and all unit vectors $\psi \in \mathbf{C}^n$:

$$\log \|T\|_p \geq \log \|T\psi\|_q - \log \|\psi\|_p,$$

where $\frac{1}{p} + \frac{1}{q} = 1$. Equality holds here for $p = 2$, hence we may differentiate with respect to $\frac{1}{p}$ at $\frac{1}{p} = \frac{1}{2}$:

$$f'_T(\frac{1}{2}) \geq -h(T\psi) - h(\psi).$$

It follows that $h(T\psi) + h(\psi) \geq -\log m$. □

The equality (2) is optimal if $|T_{ij}| = 1$ for some pair (i, j) and in the case of complementary observables, when

$$T_{jk} = \frac{1}{\sqrt{n}} e^{\frac{2\pi ijk}{n}}.$$

In general however, f_T will be strictly convex, so that $f'_T(\frac{1}{2}) < \log m$ and no ψ exists reaching equality in (2).

REFERENCES

- [BBM] I. Białyński-Birula, J. Mycielsky: "Uncertainty Relations for Information Entropy in Wave Mechanics"; *Commun. math. Phys.* 44 (1975) 129-132.
- [Deu] D. Deutsch: "Uncertainty in quantum measurements" *Phys. Rev. Lett.* 50 (1983) 631-633.
- [HLP] G. Hardy, J.E. Littlewood, G. Pólya: *Inequalities*. Cambridge University Press 1934.
- [Kra] K. Kraus: "Complementary observables and uncertainty relations"; *Phys. Rev. D* 35 (1987) 3070-3075.

[**MaU**] H. Maassen, J. Uffink: "Generalized Entropic Uncertainty Relations"; *Phys. Rev. Lett.* *60* (1988) 1103-1106.

[**Rie**] M. Riesz: "Sur les maxima des formes bilinéaires et sur les fonctionnelles linéaires"; *Acta Math.* *49* (1927) 465-497.